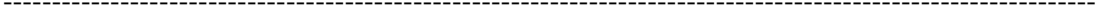


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**Impact of three-dimensional sphericity and roundness
on heat transfer in granular materials**

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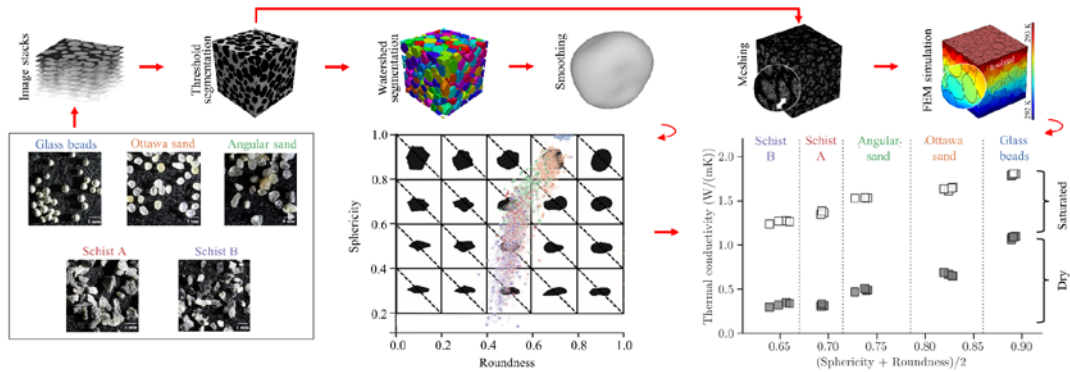
41 **Abstract**

42 Knowledge of particle morphology is vital to understand the behaviour of natural
43 geomaterials including heat transfer. The effects of particle shape on heat transfer have been
44 mostly quantified with two-dimensional (2D) particle descriptors or at most with a single three-
45 dimensional (3D) descriptor. However, these particle shape descriptors may fail to capture the
46 shape of all irregular particles. To redress this issue, we developed a method to reconstruct
47 particles from micro-computed tomographic (μ CT) images and to extract 3D sphericity and
48 roundness of individual particles in the assembly. Sphericity and roundness of five real sand
49 packings are calculated using the new proposed method. Furthermore, the effective thermal
50 conductivity (ETC) of each sample is estimated using finite element modelling. Our results
51 show that packings with higher sphericity or roundness tend to render higher ETC. A further
52 examination of the microstructure in the assemblies indicates that sphericity or roundness
53 corresponds to inter-particle contacts.

54 **Keywords:** Heat transfer, Particle shape, Granular material, Particle reconstruction.

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Graphic abstract



Five granular materials of similar particle size but different shape are scanned. From the resulting micro-CT images, the particle shape of each individual particle is automatically computed using an algorithm introduced here that comprises a series of steps shown on the top row, resulting in now well-defined Sphericity and Roundness shown in the bottom middle figure. The thermal conductivity of each material is calculated using particle scale FEM to finally unveil the relationship between particle shape and thermal conductivity, as shown in the bottom right.

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62 1 Introduction

63 Knowledge of effective thermal conductivity (λ_{eff}) of granular systems, for example, is
64 required in design, prediction and control of processes in many engineering applications such
65 as radioactive waste disposal [1], geological carbon dioxide storage [2], hydrocarbon energy [3]
66 and geothermal engineering [4]. As the λ_{eff} is affected by the microstructure [5, 6] in the
67 granular materials, to quantify the impact of microstructural features on λ_{eff} enhances the
68 fundamental understanding of heat transfer processes. Particle shape is a fundamental feature
69 in characterising the microstructure of natural soils and rocks. Hence, it is vital to fundamentally
70 understand how particle shape affects the λ_{eff} of these geomaterials and of other granular
71 materials.

72 Many experiments and simulations show that particle shape contributes significantly to the
73 mechanical [7-10], hydraulic [11-14] and thermal [15-17] behaviour of geomaterials. Even
74 though the correlation between particle shape and λ_{eff} has been quantified by introducing two-
75 dimensional (2D) particle descriptors [15, 18], the easily measured 2D particle shape
76 descriptors [10, 19-21] (e.g., circularity, sphericity, roundness, aspect ratio, convexity,
77 compactness and solidity) can be random and inaccurate as they may depend on the projection
78 direction. This is particularly the case for irregular and platy particles. The relationship between
79 particle shape and thermal conductivity were also studied by using three-dimensional (3D)
80 sphericity to characterise the particle shape [22, 23]. However, (synthetic) particles with only a
81 certain shape such as sphere and cylinder were selected for those studies. Moreover, the 3D
82 sphericity defined [24] as the ratio of the surface area of the equivalent sphere of a particle to
83 the real surface area of the particle can be used to represent the particle shape. However, this
84 sphericity definition cannot distinguish, for example, between particles with a disc-shape and
85 particles with a rod-shape. Discrete element method (DEM) has also been used to generate
86 ellipsoids varying in aspect ratio to study the effect of particle shape on λ_{eff} [17]. However,
87 the artificial ellipsoid or super-ellipsoid [25] from DEM are different from many real natural
88 sands. As other definitions of sphericity [24, 26-29] and other 3D particle shape descriptors
89 such as roundness [24] are available, these 3D particle shape descriptors (and definitions)
90 should be identified and chosen carefully to ensure they can properly describe *all* possible
91 irregular particles in a packing or assembly.

92 Particle shape and particle arrangement (i.e., microstructure) influence the behaviour of the
93 assemblies. To rigorously study their impact on the thermal (and other) response of
94 geomaterials, one would prefer to derive shape descriptors from the same sample (and of the
95 particle assembly) for testing thermal parameters. Since particle shape descriptors are
96 calculated based on the particle geometry, individual particles should be extracted from the
97 particle assembly without disturbing the structure of the sample rather than selecting a few

98 representative particles from the sample. Techniques such as resin impregnation and sequential
99 2D (physical) slicing and polishing of surfaces throughout the sample and the use of common
100 microscopic facilities are effective but time and labour intensive, and destroy the samples [20,
101 30, 31]. Alternatively, X-ray computed tomography (CT) is one of the non-destructive image
102 techniques able to visualise 3D particles in the samples [32-34]. Even though CT techniques
103 have been used in the study of mechanical response [35, 36] and fluid flow [37] in granular
104 materials, they have rarely been applied to the study of heat transfer. CT techniques can
105 generate sequential cross-sectional images of the geomaterial at a certain interval (image
106 resolution) and the solid phase can be detected based on the greyscale of each image. From the
107 stacking of the solid phase in all images, one can construct a “voxelated” geometry. However,
108 the particles in the geometry may overestimate [38] the inter-particle contact area and even
109 appear or cemented to each other. As the contact area is related to the thermal conductance [39],
110 post-processing image techniques are required to identify “true” individual particles.

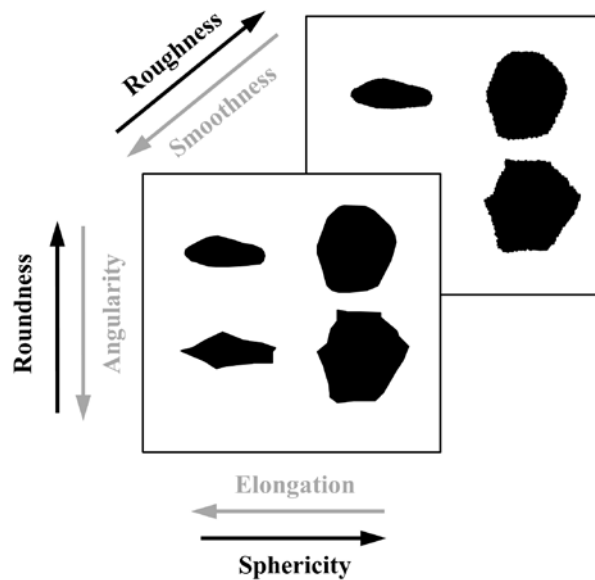
111 Based on the collection of individual particles extracted from voxel-constructed geometry,
112 the particle surface area can be measured using the boundary voxels, and the particle volume
113 can be computed based on the number of voxels inside each particle [38, 40]. However, the
114 voxelated particle surface has a saw-tooth pattern and overlapping voxels, leading to an
115 overestimation [38] which means the computed value may further affect the accuracy of
116 measuring sphericity and roundness. Moreover, since 3D roundness are widely accepted [24,
117 27, 41] as the ratio of the average radius of the particle corners and the radius of the maximum
118 inscribed circle, the voxel-constructed particle surface cannot be used to calculate the correct
119 curvature for identifying the corners. Consequently, smoothing the voxel-constructed particles
120 before calculating the 3D sphericity and roundness is required. These issues will be addressed
121 in the present study.

122 The study aims to investigate the effect of particle shape on the effective thermal
123 conductivity of granular materials. To achieve this, various particle shape descriptors are
124 compared to then determine a combination of shape descriptors that can actually cover particle
125 with a wide range of different shapes. Moreover, the theories used to compute the selected
126 particle shape descriptors are explained. A framework is introduced to generate individual
127 smooth particles based on CT images and the methodology is exemplified using five real sand
128 packings. The experimental and numerical methods used to measure the effective thermal
129 conductivity are explained in detail, followed by quantitative analysis on how and why particle
130 shape and microstructure affect heat transfer in granular materials.

131 **2 Particle shape descriptors**

132 Particle shape can be generally described at three different length scales [7] (Fig. 1):
133 Sphericity/Elongation, Roundness/Angularity and Smoothness/Roughness. At the scale of the

134 equivalent particle diameter, sphericity indicates the global form of the particle and describes
 135 the proximity of a particle to a sphere while elongation describes the opposite trend. At
 136 approximately one tenth of the scale of the equivalent particle diameter, roundness characterises
 137 the particle shape at a smaller scale than the scale of sphericity. Roundness is a local feature
 138 and represents the extent to which the corners and edges of the particle have been rounded, its
 139 counterparty is known as angularity. Smoothness or roughness describes the smallest scale and
 140 indicates the plateness of the particle surface. Since images with resolution higher than 1 μm
 141 [42] are typically required to calculate the roughness, the present work focuses on the
 142 implementation of sphericity and roundness because the resolution of CT scans is typically
 143 lower. The CT scanned images used in this work, for example, have a resolution of
 144 approximately 13 μm .



145
 146 Fig. 1 Particle shape descriptors at three different scales: Sphericity, Roundness and
 147 Smoothness.

148 *2.1 Three-dimensional sphericity and roundness*

149 A number of definitions (and corresponding equations) have been proposed in the literature
 150 to calculate three-dimensional (3D) sphericity and roundness. Table 1 summarises some of
 151 these definitions and formulae to compute 3D sphericity (S1 – S5) employing different
 152 parameters. While particle volume V and surface area SA are used in S1 and S2, the radius of
 153 maximum inscribed and minimum circum-scribed sphere ($r_{\text{max-in}}$ and $r_{\text{min-cir}}$, respectively) are
 154 adopted in S3. The length of principal axes of the particle d_i is employed in S4 and S5. However,
 155 none of these descriptors can distinguish all particles with different shapes. For example, S1
 156 and S2 cannot distinguish between disc-shape particles and rod-shape particles because they
 157 may have the same surface area and volume, as depicted in Fig. 2(a). In Fig. 2(b) and Fig. 2(c)
 158 S3 cannot recognise the difference between particle 1 and particle 2 because they have the same
 159 maximum inscribed circle and the minimum circum-scribed circle. Moreover, S4 and S5 cannot

160 distinguish particle 1 and particle 2 in Fig. 2(c) since they have the same principle axes if
 161 considering the two particles have the same thickness in 3D.

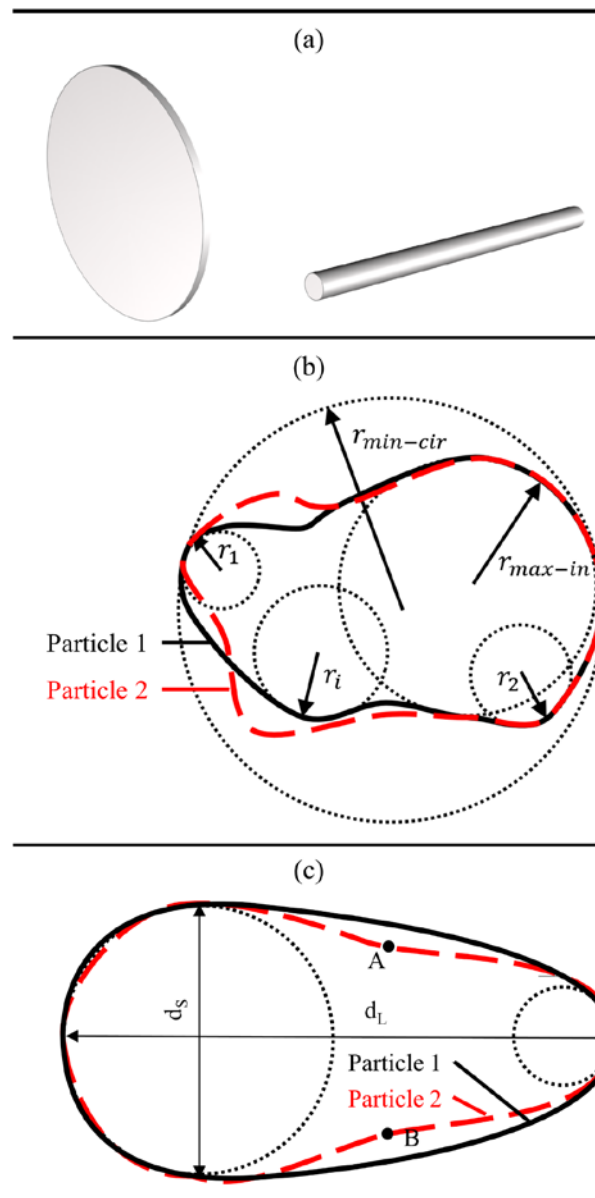
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Table 1 A summary of various definitions of sphericity

Notation	Equation	Description	Reference
S1	$\frac{\sqrt[3]{36\pi V^2}}{SA}$	The ratio of the surface area of the equivalent sphere of a particle to the real surface area of the particle. V is particle volume and SA is particle surface area.	[24]
S2	$\frac{36\pi V^2}{SA^3}$	The cubic order of S1.	[26]
S3	$\frac{r_{\max\text{-in}}}{r_{\min\text{-cir}}}$	The ratio of the radius $r_{\max\text{-in}}$ of the maximum inscribed circle (sphere in 3D) to the radius ($r_{\min\text{-cir}}$) of the minimum circum-scribed circle of a particle shown in Fig. 2(b).	[27]
S4	$\frac{d_S}{d_L}$	The ratio of the maximum axial length (d_S) to the minimum axial length (d_L) of a particle.	[28]
S5	$\sqrt[3]{\frac{d_S^2}{d_L d_I}}$	Fitting a particle to an ellipsoid, d_S , d_I and d_L are the shortest, intermediate and the longest axial length of the fitted ellipsoid.	[29]

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Fig. 2 Examples of potential shortcomings of the various definitions of sphericity and roundness: (a) Example of two particles of different shape but same sphericity (S1 and S2 in Table 1); (b) Example of two particles of different shape but same sphericity (S3 in Table 1) and roundness. (c) Example of two particles of different shape but same sphericity (S4 and S5 in Table 1) and roundness

In contrast to the various definitions of sphericity, only few equations have been proposed to compute 3D roundness. Equation 1 below is used to calculate 3D roundness as the ratio of the average radius of the particle corners to the radius of the maximum inscribed sphere (see Fig. 2(b)):

$$R = \frac{\sum r_i/N}{r_{max-in}} \quad (1)$$

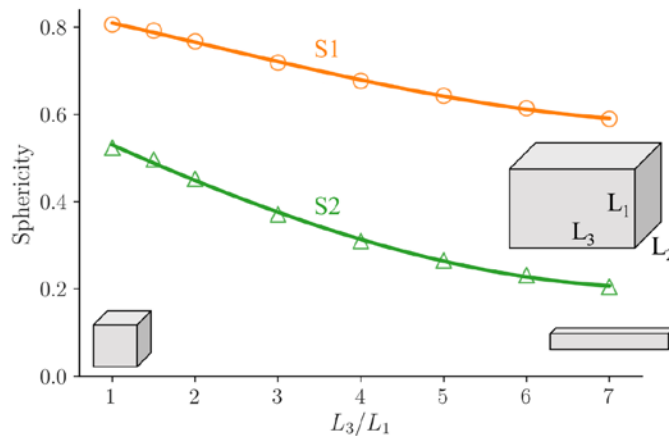
177 where r_i is the radius of corners in a particle as shown in Fig. 2(b), N is the total number of
 178 corners and r_{max-in} is the maximum inscribed circle of a particle. However, the roundness of
 179 the particles in Fig. 2(c) cannot be distinguished if neglecting the local feature on particle
 180 surface where the radius of curvature is larger than the maximum inscribed circle.

181 Since neither sphericity nor roundness alone can distinguish *all* particles with different
 182 shapes, combinations of sphericity and roundness have been used to ensure that all particles
 183 with various shapes can be identified [41]. From the previous explanation, particles shown in
 184 Fig. 2(c) still cannot be distinguished even if roundness is combined with sphericities S4 or S5.
 185 Furthermore, S3 cannot distinguish the two particles in Fig. 2(c), hence using roundness with
 186 S3 still fails to recognise the two different particles. Consequently, the remaining S1 and S2
 187 definitions are good candidates which together with roundness could characterise particle shape
 188 unambiguously.

189 To analyse the performance of S1 and S2, a hexahedron with dimensions $L_1 \times L_2 \times L_3$ is
 190 adopted but keeping L_1 equal to L_2 and then changing the ratio of L_3 to L_1 . According to the
 191 results in Fig. 3, a hexahedron with L_3/L_1 varying between 1 and 7, shows a sphericity that
 192 shifts from 0.81 to 0.59 when using definition S1, while it changes from 0.52 to 0.21 when
 193 using definition S2. Since sphericity is conceptually defined as the proximity of a particle shape
 194 to a sphere (sphericity is 1) and a hexahedron with $L_3/L_1 = 7$ is far from a sphere, assigning
 195 0.59 to the sphericity of such an elongated particle may be intrinsically unacceptable, as one
 196 would expect a much lower value.

197 As a consequent of the above discussions, in this work sphericity as defined in S2 in Table
 198 1 and roundness as defined in Equation 5 are selected as the combination of descriptors to better
 199 characterise the shape of particles in natural granular materials like sands.

200



201

202 Fig. 3 Comparison of sphericity S1 and S2 defined in Table 1 of a hexahedron with

203 varying dimensions ($L_1=L_2$, $1 \leq L_3/L_1 \leq 7$)

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206 2.2 Calculation of sphericity and roundness

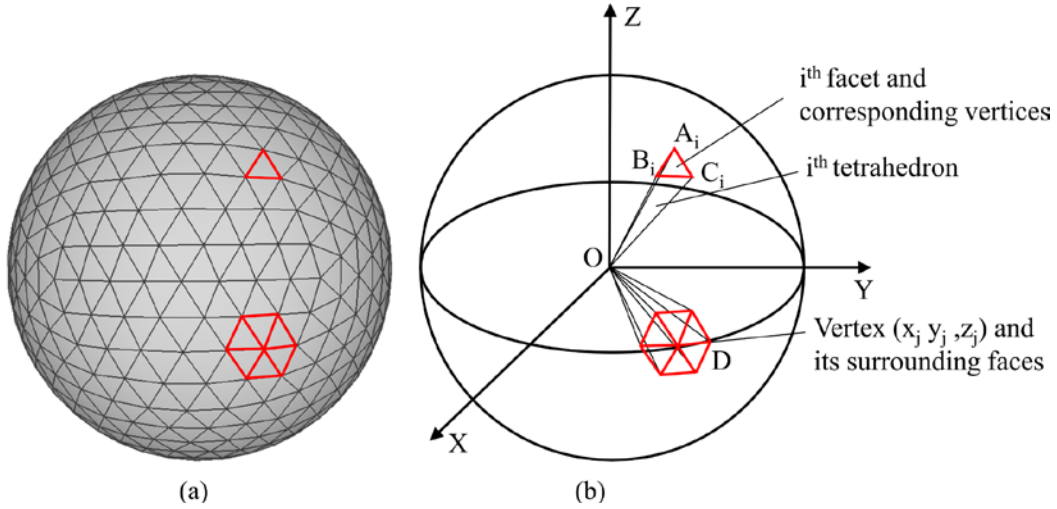
207 From the last section, it is known that the particle surface area and volume are required to
208 compute sphericity whereas curvature is required for calculating roundness. While Fiji [43] is
209 available to compute the particle surface area and volume, the computations are based on voxel-
210 constructed particles which have saw-tooth boundary patterns. The saw-tooth boundary
211 patterns may lead to overestimation of real particle surface area and an incorrect calculation of
212 curvature, the computation of sphericity and roundness should be based on smooth particles.
213 However, no commercial software or open-source code is available to achieve smooth particle
214 geometry and calculate 3D sphericity and roundness using the equations we selected. In this
215 work, we first identify each individual particle and then the smoothing particle surface is made
216 up of triangle surface meshes (similar to [41] as shown in Fig. 4(a). Therefore, the surface area
217 of each facet on the particle surface is easily computed, and the summation of all the triangular
218 facet areas is the surface area of the particle (Fig. 4(b)). To compute the individual particle
219 volume, the following procedure is followed. A point inside the particle together with the three
220 points of a single triangular facet form a tetrahedron, so the total volume of the tetrahedrons
221 related to all facets is the particle volume (Fig. 4(b)). Thus, to calculate the particle surface area
222 and volume, Equations 2 and 3 are used. Given particle volume and surface area, then S2 can
223 be calculated as the sphericity of each particle.

$$SA = \sum_{i=1}^n \frac{1}{2} |\overrightarrow{A_i B_i} \times \overrightarrow{B_i C_i}| \quad (2)$$

$$V = \sum_{i=1}^n \frac{1}{6} |\overrightarrow{OA_i} \cdot (\overrightarrow{OB_i} \times \overrightarrow{OC_i})| \quad (3)$$

224 where O is the particle centre calculated as the average coordinates of all vertices in the particle,
225 A_i , B_i and C_i are three vertices of the i^{th} triangular facet and n is the total number of the facets
226 on the particle surface.

227



228 (a) 229 Fig. 4 A particle surface is made up of triangular meshes. (a) Triangular meshes of a 230 sphere, (b) A diagram to illustrate the principles of computing particle surface area, volume 231 and curvature at vertices (after [41]). 232

233 To calculate the roundness of a particle, the first step is to identify corners in the particle. 234 An available criterion is that considering a vertex as a corner when its radius of curvature is 235 smaller than the maximum inscribed sphere. The principal curvatures (the maximum curvature 236 κ_{\max} and minimum curvature κ_{\min}) are two common measurements [44], and quadratic fitting 237 in MeshLab [45] can be used to compute them at all vertices. The key of this method is to use 238 rings of a vertex (e.g., the first ring of vertex D in Fig. 4(b) is composed of the six vertices 239 around it) to fit a micro-surface by generating a quadratic polynomial equation. Based on the 240 equation, a Hessian matrix [44] can be established and the eigenvalues of the matrix are the 241 principal curvatures. Subsequently, the principal curvature radii are computed as $r_{\min} =$ 242 $1/\kappa_{\max}$ and $r_{\max} = 1/\kappa_{\min}$. Between the two radii, r_{\min} is more reasonable to be used for 243 identifying corners because more local features on the particle surface are considered. Then, 244 the corners in a particle can be found using Equation 4.

$$g(k) = \begin{cases} 1 & \text{if } r_{\min} < r_{\max-in} \\ 0 & \text{if } r_{\min} \geq r_{\max-in} \end{cases} \quad (4)$$

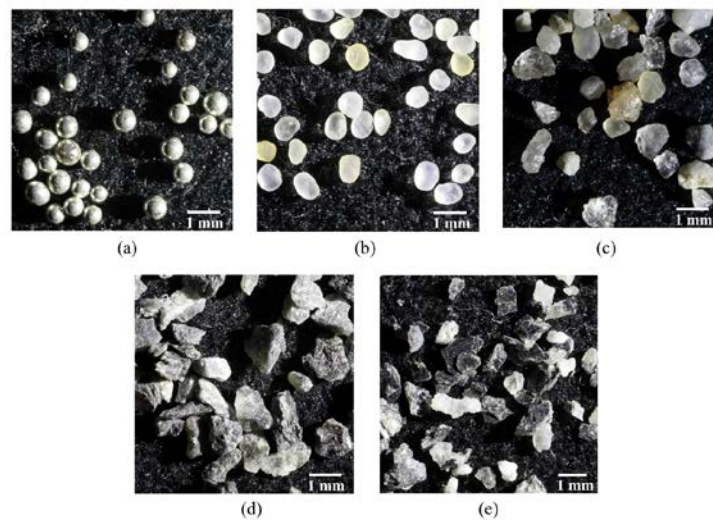
245 where 1 indicates that the vertex is a corner, while 0 is not, $r_{\max-in}$ is the radius of the 246 maximum inscribed sphere of the particle. By introducing $g(k)$ into equation 5, the 3D 247 roundness is computed as:

$$R = \frac{\sum g(k)r_{\min} / N}{r_{\max-in}} \quad (5)$$

248 3 Particle reconstruction

249 3.1 Granular material samples

250 To test the framework and to analyse the impact of particle shape on effective thermal
 251 conductivity, five sands with increasing irregularity of particles are selected (Fig. 5): (a) glass
 252 beads are near-spherical particles made from silica; (b) Ottawa sand is sieved to pass No.20
 253 (850 μ m) mesh and be retained on No.30 (600 μ m), the achieved 20-30 standard sands are in
 254 line with ASTM standard C778 [46]. These particles consist of quartz [47] and are rounded
 255 during long-term erosion. (c) an Angular sand also containing a high proportion of quartz with
 256 more irregularly shaped particles when compared to the Ottawa sand. (d) Crushed schist sand
 257 A is even more irregular than the Angular sand and the particles in Crushed schist A are made
 258 of chlorites. (e) Crushed schist B have particles with the most complex shape and more than
 259 half of the particles are platy and elongated [48] because they consist of quartz and biotite [49].
 260 The Crushed schist rock B is made from a particular metamorphic rock collected from the
 261 Delamarian Fold Belt in western Victoria, Australia. The particle size of each sample is shown
 262 in Table 2. Crushed schist rock A and schist rock B are prepared to have the same mean particle
 263 size and range in the laboratory.
 264



265
 266 Fig. 5 Micro-pictures of five granular materials with particles of different shape. (a) Glass
 267 beads, (b) Ottawa sand, (c) Angular sand, (d) Crushed schist rock A and (e) Crushed schist
 268 rock B.

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 270

Table 2 Particle size of studied granular materials

Sample	d_{50} (mm) *	d_{50} (mm) ^	Particle size range (mm) *	Particle size range (mm) ^
Glass beads	0.60	0.60	0.50 – 0.70	0.40 - 0.80
Ottawa sand	0.73	0.76	0.60 – 0.85	0.58 - 0.94
Angular sand	0.89	0.68	0.60 – 1.18	0.39 – 0.99
Crushed schist rock A	0.84	0.58	0.50 – 1.18	0.23 – 0.95

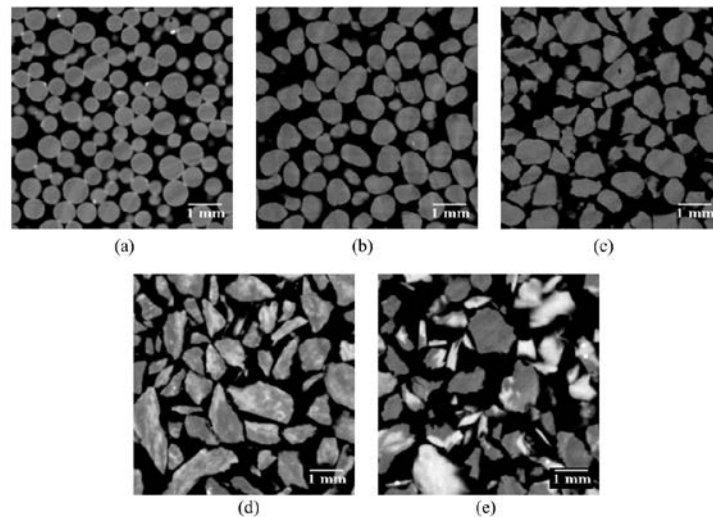
Crushed schist rock B	0.84	0.61	0.50 – 1.18	0.16 - 1.10
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271 * Diameter from sieve analysis;

272 ^ Equivalent diameter calculated from particle volume using CT reconstruction. Then box plots are
 273 drawn, then d_{50} is the median value and particle size range is the range in the box by ignoring the outliers.
 274

275 These five materials were used to prepare samples in a 25 mm diameter and 25 mm height
 276 cylindrical containers made of aluminium using the air-pluviation technique to ensure the
 277 homogeneity of the samples and density similarity. The samples are scanned using x-ray
 278 computed tomography (CT) to generate sequential grey-scale images with a resolution of 0.013
 279 mm. This resolution lies between the recommendations by [Wiebicke, Andò, Herle and Viggiani](#)
 280 [50]. Their work suggested that 10-15 μm per voxel and more than 10 pixels across each grain
 281 diameter are required to capture the interparticle contact with improved image processing
 282 techniques. The greyscale in each voxel represents the image intensity which relates to the
 283 density of the scanned materials. Typical cross-sectional image of each sample is shown in Fig.
 284 6. Particles in Glass beads, Ottawa sand and Angular sand samples have similar grey scale even
 285 though some defects exist on the boundary of glass beads because of their coating. In contrast,
 286 particles in crushed schist rock samples have significant different greyscale because of their
 287 complex mineral components.

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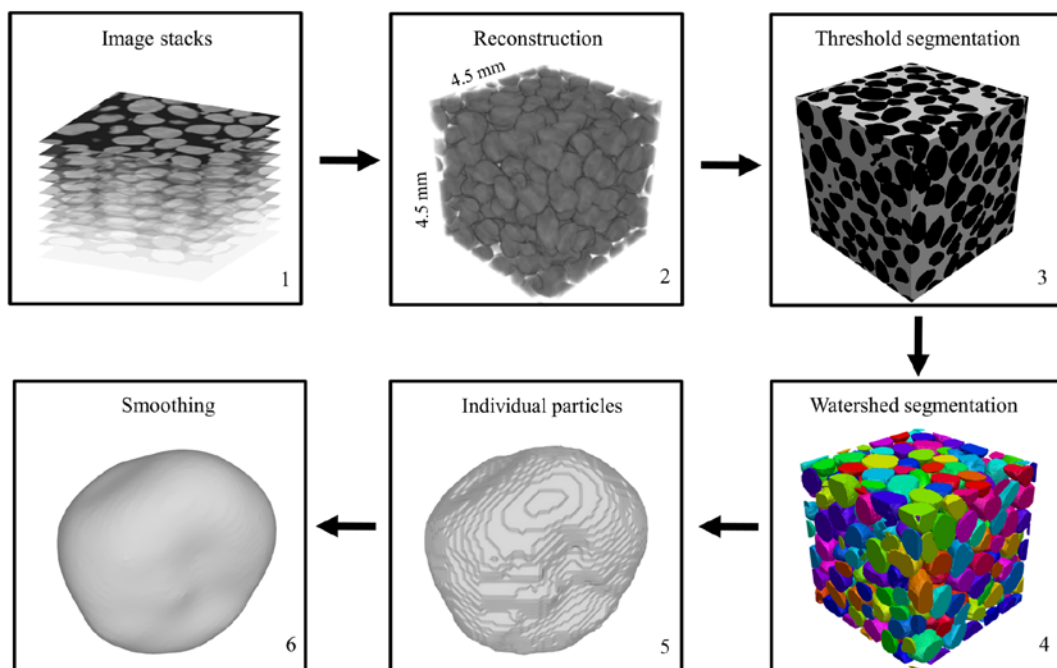
290 Fig. 6 . CT scanned images of the five granular material packings. (a) Glass beads, (b)
 291 Ottawa sand, (c) Angular sand, (d) Crushed schist rock A and (e) Crushed schist rock B

292 3.2 Framework

293 The geometries of *individual* particles can be acquired using the X-ray CT scans of separate
 294 particles placed in a container filled with a (low density) silicon filler [41]. However,
 295 identifying individual particles in the samples (i.e., assembly of particles) is not always
 296 straightforward. Hence, segmentation is required to identify and extract information of

297 individual particles from the specimen. To accurately calculate 3D sphericity and roundness, a
 298 smooth digital particle geometry with the “same” surface area and volume as the real particle
 299 is required. Here we combined different techniques and propose a series of steps to reach this
 300 goal. This section introduces a framework (Fig. 7) to reconstruct individual smooth particles
 301 from a stack of cropped X-ray CT scanned images (Step 1). An open source image processing
 302 software Fiji [43] is used to reconstruct the sample micro-geometry (Step 2) and to execute
 303 threshold segmentation to identify the solid phase of the particle assemblies (Step 3). [Iassonov,](#)
 304 [Gebrenegus and Tuller \[51\]](#) applied various threshold segmentation techniques (e.g. local
 305 thresholding, global thresholding, region growing methods, deformable surface, probabilistic
 306 clustering and Bayesian methods) to porous media. In their work, they found that *local*
 307 thresholding can generate satisfactory results but are sensitive to the initial input, while the two
 308 *global* threshold segmentation methods (Otsu and the Rilder’s [50, 52-54]) produced
 309 satisfactory results without this limitation [38, 50]. Therefore, in our work, the multilevel Otsu
 310 threshold segmentation method is applied to the greyscale images. Then a watershed
 311 segmentation with 6-connectivity [38, 55] in MorphoLibJ library [26] is adopted to split
 312 (“artificially”) connected particles based on the threshold segmentation geometry (Step 3). The
 313 6-connectivity is selected herein because it shows to reduce the overestimation of particle
 314 surface area and volume [38]. Boundary particles are discarded because they are part of integral
 315 particles resulting from generating the cubic subsamples in Step 1. The individual voxel-
 316 constructed particles are thus obtained (Step 5) and then the Taubin smoothing algorithm [56]
 317 is used to smooth each particle (Step 6).

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Fig. 7 A framework is used to reconstruct individual smoothing particles

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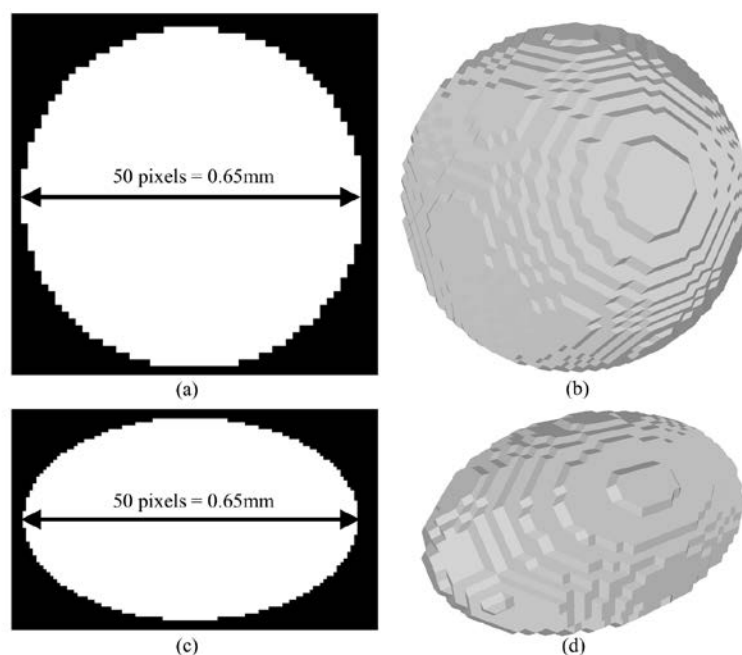
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324 3.3 Smoothing method

325 In Step 6 of the aforementioned framework, a smoothing algorithm is required to polish the
326 saw-tooth surface of voxelated particles and preserve the general particle morphology.
327 Common iterative mesh smoothing algorithms utilise “processing mesh normals” scheme [57-
328 59]. Most traditional smoothing algorithms, including Gaussian smooth and Laplacian smooth,
329 have a general shrinkage problem which means a particle may collapse to a point when
330 smoothing is implemented for a large number of iterations. Alternatively, the Taubin λ/μ
331 method [56] consists of two Gaussian smoothing algorithms in each step: λ is positive and
332 relates to the shrinkage whereas μ is negative and controls the infatuation.

333 To identify the proper range of values of input parameters in the Taubin method for our
334 samples, synthetic pixelated image stacks of a sphere and an ellipsoid are generated using the
335 image analysis software Fiji [43], the resolution of these images is set at 0.013 mm, which is
336 the same as the resolution used in the CT of the tested samples. The diameter of the sphere is
337 set as 0.65 mm (50 voxels), and the three-principal axial lengths of the ellipsoid are set at 0.65
338 mm (50 voxels), 0.39 mm (30 voxels) and 0.26 mm (20 voxels). The resulting cross-section
339 images and 3D voxelated particles are shown in Fig. 8.

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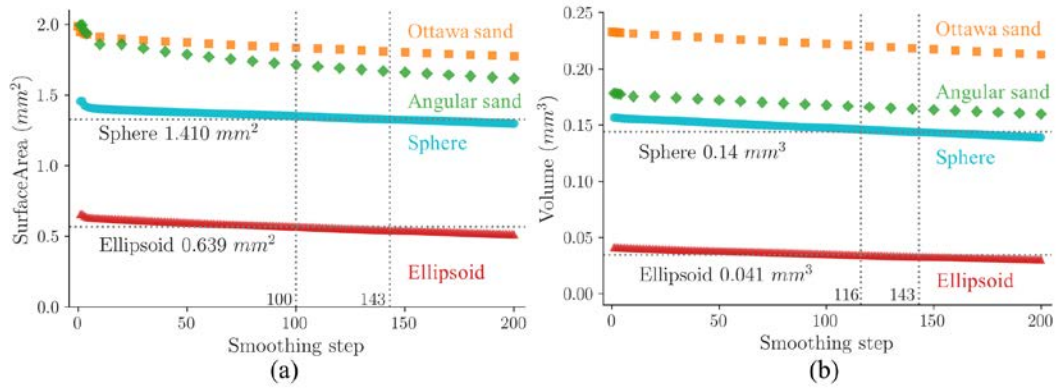
342 Fig. 8 Voxelated sphere and ellipsoid are constructed from images stacks generated from
343 Fiji: (a) 2D image slice of a sphere, (b) and the resulting voxelated sphere; (c) 2D image slice

344 of an ellipsoid and (d) the resulting voxelated ellipsoid. The resolution of images is 0.013
345 mm, same as the resolution of CT scanned images later used in his work

346

347 The Taubin smoothing method effectively generates a *low pass filter effect* that is controlled
348 by λ and μ , and the amount of attenuation is then determined by the number of smoothing steps
349 set in the algorithms. The Taubin smoothing method can preserve the geometry of the
350 reconstructed particles from CT images when $\lambda < -\mu$ [56]. Hence, the shrinkage problem in
351 Gaussian smooth and Laplacian smooth is addressed in the Taubin smoothing method.
352 However, if the surface area and volume from the voxelated particles in the initial step are
353 overestimated to begin with (typical of microCT images), these overestimations are also
354 preserved [38]. To mitigate the initial overestimation of surface area and volume of a particle,
355 $\lambda > -\mu$ is adopted in this work by setting λ as 0.6 and μ as -0.3, which also happen to preserve
356 the geometry of the particles. Together with the synthetic sphere and ellipsoid generated above,
357 two microCT imaged particles, one from the Ottawa sand and one from the Angular sand
358 samples, are also used to investigate the variation of surface area and volume with the number
359 of smoothing steps. Fig. 9 shows that the surface area gradually decreases with smoothing steps,
360 and that the surface area of the voxelated sphere reaches the theoretical value (1.410 mm^2) at
361 the 143rd smoothing step while that of the ellipsoid is reached earlier at the 100th step. As for
362 the reduction of volume, the four particles share a similar trend but at a slower rate of change
363 than that of surface area, the volume of the voxelated sphere reaches the theoretical value (0.41
364 mm^3) at the 143rd step, which is the same as that of surface area, while the ellipsoid, at the 116th
365 step. In addition to the above analysis to define the range of smoothing steps values that may
366 be appropriate for our studies, sphericity and roundness of Ottawa sand and Angular sand
367 particles are computed and shown in Fig. 10 together with their images at 0, 100 and 200
368 smoothing steps. It can be seen that the local features of the surface are preserved during the
369 process of Taubin smoothing. Fig. 10 shows that sphericity and roundness converge at early
370 smoothing steps, which is far earlier than the 143rd step found in the previous analysis. To be
371 conservative and to allow a good balance between computational time and accuracy of results,
372 143 steps (together with λ as 0.6 and μ as -0.3) are selected as input parameters in the
373 smoothing method used herein in the following calculations for all individual particles in
374 natural sands scanned in the micro-CT.

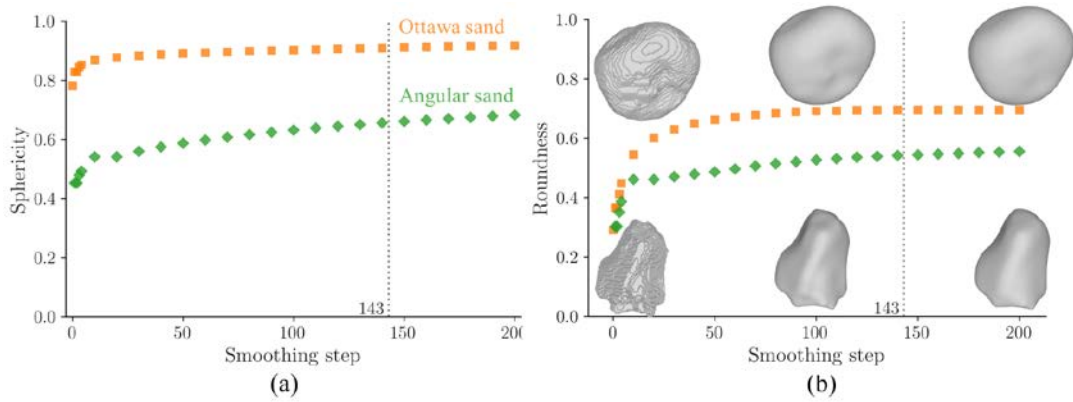
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377 Fig. 9 The surface area (a) and volume (b) decrease gradually with the increase in the number
 378 of smoothing steps for $\lambda > -\mu$ in the Taubin smoothing method.

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381 Fig. 10 The sphericity and roundness of two particles from the Ottawa sand and the
 382 Angular sand samples as a function of smoothing steps: (a) Sphericity and (b) Roundness.

383 One can also visualize the Ottawa and Angular sand particles at 0, 100 and 200 smoothing
 384 steps for comparison

385 4 Finite element simulation and measurements

386 The mineral component of a sand can affect its effective thermal conductivity, especially
 387 the content of quartz in dry samples [60]. To eliminate this effect, finite element simulation is
 388 used in this paper by assigning the same thermal conductivity to the solids of all samples.
 389 Experimental measurements, however, are used to validate the numerical results.

390 4.1 Finite element simulation

391 The geometry of a sand sample can be reconstructed from scanned CT images and used to
 392 simulate heat transfer in a similar fashion as in [61, 62] by numerically solving an elliptic partial
 393 differential equation (Equation 6), Fourier's law (Equation 7) and a continuity equation
 394 (Equation 8), using COMSOL Multiphysics [63]. The thermal conductivity of the packing are
 395 obtained by integrating the heat flux at the top and bottom boundaries using Equation 9, and
 396 their average is taken as the effective thermal conductivity of the entire sample.

$$\rho C \frac{\partial T}{\partial t} + \rho C u \cdot \nabla T = \nabla \cdot (\lambda \nabla T) \quad (6)$$

397 where, for each phase involved in the simulation, ρ is the density (kg/m), C is the heat capacity
 398 (J/(kg K)), T is the temperature (K), t is the time (s), u is the velocity vector (m/s), λ is the
 399 thermal conductivity (W/(m K)).

$$q = \lambda \nabla T \quad (7)$$

400

$$-n(q_s - q_p) = 0 \quad (8)$$

401 where n is the unit normal vector of the solid-pore interface, q_s and q_p are the heat fluxes in the
 402 particle and pore, respectively.

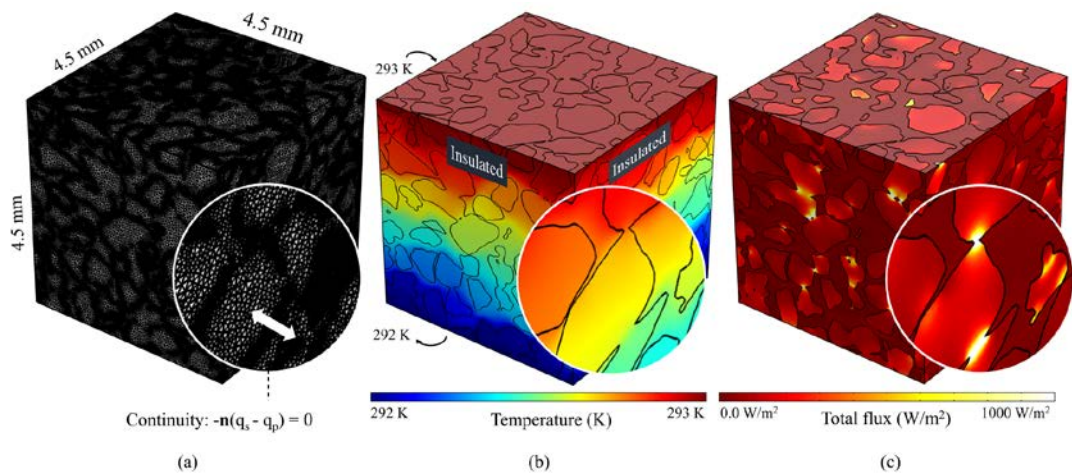
403 The effective thermal conductivity λ_{eff} ($\text{W m}^{-1}\text{K}^{-1}$) of a sample of horizontal cross-section
 404 area A (m^2) is found as:

$$\lambda_{\text{eff}} = \frac{\frac{1}{A} \int_A Q_z dA}{\frac{T_a - T_b}{L}} \quad (9)$$

405 where T_a and T_b are the prescribed temperatures at the inlet and outlet boundaries, L (m) is the
 406 height of the sample and Q_z (W/m^2) is the vertical heat flux of nodes at the inlet or outlet.

407 The mesh of the Angular sand is shown in Fig. 11(a) and it was generated in Simpleware
 408 ScanIP [64] by setting coarseness as -40 after a mesh size sensitivity analysis showed
 409 convergence to an asymptotic value of computed thermal conductivity (analysis not included
 410 here). The thermal conductivity of minerals has been suggested to be set between 1 and 8
 411 $\text{W}/(\text{m K})$ in the literature [61]. Since the aim of this study is on the effect of particle shape on
 412 effective thermal conductivity, 3 $\text{W}/(\text{m K})$ [39, 65, 66] is assigned to solid particles in all
 413 samples to mitigate the potential effects of different mineralogy on the effective thermal
 414 conductivity. Thermal conductivity of 0.025 $\text{W}/(\text{m K})$ [67] and 0.591 $\text{W}/(\text{m K})$ [67] is
 415 assigned to the void space within the packings for simulating heat transfer in dry and water-
 416 saturated granular materials, respectively¹. The boundary temperature on the top surface is
 417 prescribed at 293 K, while at 292 K on the bottom surface to generate a small thermal gradient,
 418 other boundaries are considered as insulation as shown in Fig. 11(b). With this material and
 419 boundary conditions, the system is numerically solved for temperature distribution and heat
 420 fluxes are estimated to then derive an effective thermal conductivity as described in detail in
 421 [61].

¹ Thermal conductivity of bound water may differ from that of free water, while this effect is negligible in coarse grained packings, it may need to be considered in fine grain packings (e.g., clays).



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Fig. 11 Finite element mesh and heat transfer simulation results of the Angular sand packing (a) Mesh (b) Temperature distribution and (c) Total heat flux distribution

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426 4.2 Experimental measurements

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In order to validate finite element simulating results, the effective thermal conductivity of Ottawa sand, Angular sand and Crushed schist A are measured in the laboratory. Samples are also air-pluviated into a PVC cylindrical container with a diameter of 50 mm and a height of 120 mm to achieve similar homogeneity as shown in the CT images (Fig. 6). The measurements use a 100 mm thermal needle probe (KD2 Pro thermal properties analyser from Decagon Devices, Inc) following ASTM standard D5334-14 [68]. The thermal needle has a diameter of 2.4 mm which is larger than the largest particle diameter, which results in a good accuracy of the measurement at $\pm 10\%$ for 0.2 – 0.4 W/(mK) [69].

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5 Results

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5.1 Finite element simulating validation

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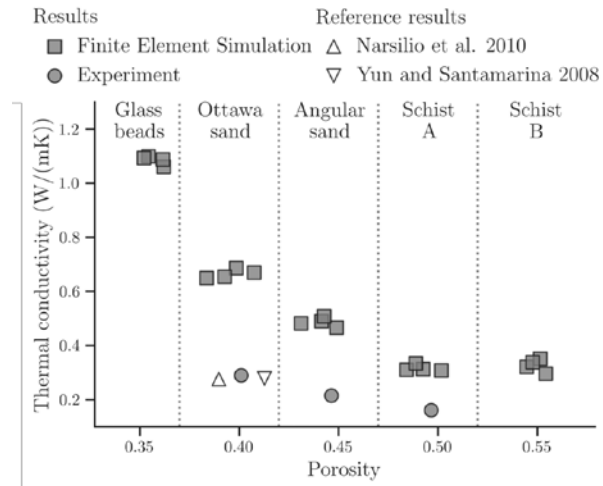
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Four subsamples from each scanned natural sand are extracted to calculate the effective thermal conductivity using the finite element simulation. Fig. 12 shows that the thermal conductivity decreases with porosity, and this trend agrees with the results presented previously by Yun and Santamarina (2008) [65] and Narsilio et al (2010) [61]. Fig. 12 also reveals that the numerical values are higher than that from experimental measurements but in an acceptable range [61]. The overestimation may result from the contact condition in both physical testing and numerical simulation. The needle probe testing has an accuracy of $\pm 10\%$ [69] and the error of finite element simulation come from its limitation in capturing particle surface roughness and particle geometry irregularity [61].



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452 *5.2 Effect of particle shape on thermal conductivity*

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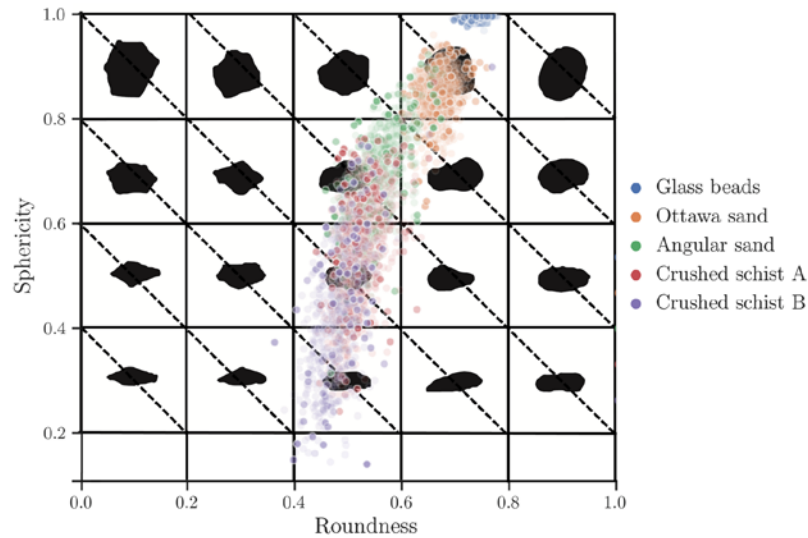
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Fig. 12 The effective thermal conductivity of five natural sands from finite element simulations are larger than the experimental measurements, which may be because numerical simulation does not capture particle surface roughness and particle geometry irregularity [51]

A sub-sample X-ray CT image stack from each natural sand is selected to extract smooth individual particles using the framework summarised in Fig. 7. The average number of particles in each sub-sample is 150. This number of particles considered sufficiently representative for studies of uniform sands [20]. Within each sub-sample, for each smooth particle in it, its sphericity and roundness are calculated using S1 in Table 1 and Equation 5. Readers must recall that sphericity (S) and roundness (R) characterise particle shape at two different length scales: the overall form of the particle (S) and local features (i.e., the shape of “edges or corner” in the particle, R) as it can be seen in Fig. 1. Each coloured point in Fig. 13 corresponds to the calculated sphericity and roundness of each particle in the sub-sample assembly of particles. A particle shape chart [70] is also included as a background in Fig. 13 to provide a rough visualisation of the particle geometry. From Fig. 5, we already qualitatively know that the ascending ordering of angularity of the five sands is: Glass beads, Ottawa sand, Angular sand, Crushed schist A and Crushed schist B. In Fig. 13, we can also observe this trend: irregular particles have lower sphericity and roundness than more regular (and round) particles. Moreover, the range of sphericity and roundness in irregular sands have a wider distribution.



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469 Fig. 13 Irregular particles have lower sphericity and roundness than round particles.

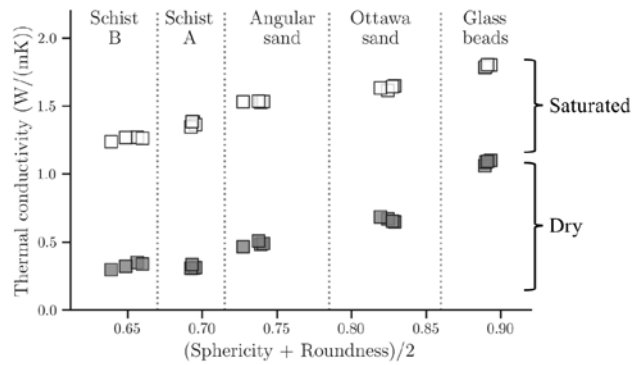
470 Packings assembled by irregular particles have a wider range of sphericity and roundness.

471 The background refers to [61]

472

473 After determining sphericity and roundness, their average can be used as a uniformed
 474 character of particle shape for the assembly of particles in each sub-sample [27]. The effect of
 475 particle shape on effective thermal conductivity can now be analysed with the quantitative
 476 particle shape descriptor. In this work, both the thermal conductivity of the five sands in dry
 477 and static water saturated conditions are computed as shown in Fig. 14 for five sub-samples of
 478 each sand. According to Fig. 14, the effective thermal conductivity increases when the average
 479 of sphericity and roundness increases. Furthermore, the trend in dry samples and water
 480 saturated samples are similar; however, the rate of increase in dry samples is slightly higher.
 481 The difference of the increasing rate may arise from the change in the main heat transfer
 482 pathways in the granular materials. As shown in Fig. 11 (c), the heat flux in *dry* granular
 483 materials is high at the particle contacts which is the main particle scale path for heat transfer
 484 in dry granular materials. However, the main heat transfer path in saturated sand is particle-
 485 fluid-particle whose contribution to the λ_{eff} is between 39.6% (in saturated Glass beads) and
 486 74.1% (in saturated Crushed schist B). Since the pore space now is filled with water, the λ_{eff}
 487 of saturated Crushed schist B is around four times larger than the λ_{eff} of dry Crushed Schist B
 488 and larger than the λ_{eff} of dry Glass beads.

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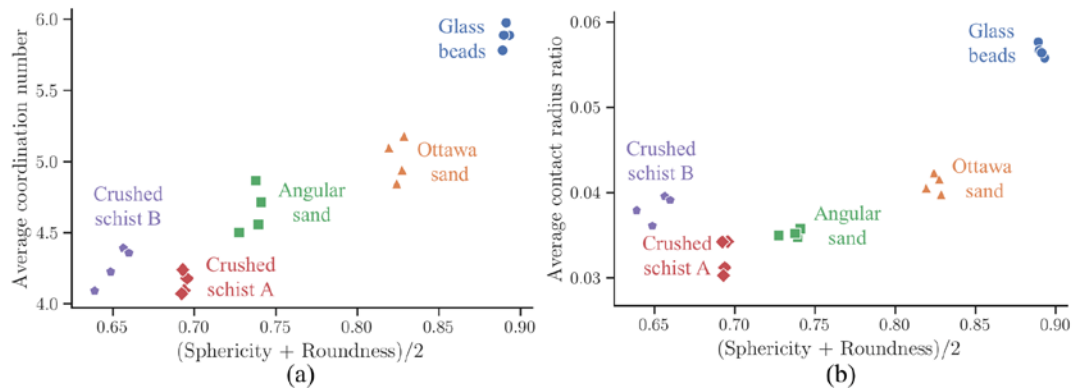
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491 Fig. 14 Effective thermal conductivity increases when the average of sphericity and
 492 roundness increases in both dry (grey) and water saturated sands (white)

493 *5.3 Effect of particle shape on microstructure*

494 Since the inter-particle contact is the main heat transfer pathway in dry granular materials,
 495 coordination number and the inter-particle contact area are two microstructure properties
 496 believed to have a significant influence on the overall thermal conductivity. Accordingly, we
 497 further study the relationship between these two microstructure parameters and particle shape
 498 descriptors. After watershed segmentation (step 4 in Fig. 7), coordination number and contact
 499 area are calculated, and then the average coordination number and the average contact radius
 500 ratio for each microCT-ed sub-samples are calculated as well. The contact radius ratio is defined
 501 as the ratio between the radius of contact area and the radius of particles. Fig. 15 shows that
 502 both coordination number and contact radius ratio increase with the increase of the average
 503 value of sphericity and roundness, except for the Crushed schist B. Particles in the Crushed
 504 schist A and the Crushed schist B samples have similar average coordination number but
 505 Crushed schist B shows a higher average contact radius ratio, which may be because almost
 506 half of the particles in the Crushed schist B sample are platy and elongated (Fig. 13) particles
 507 resulting in larger interparticle contact areas. The increase of coordination number in the
 508 samples having more regular (round) particle coincide with the increment of λ_{eff} shown in Fig.
 509 14. In general, the increase of contact ratio in more regular particle packings also coincide with
 510 the increment of λ_{eff} except for the packings containing very irregular and platy particles.
 511 These observations indicate the impact of particle shape represented by the three-dimensional
 512 sphericity and roundness on heat transfer in dry granular materials is more originated from its
 513 influence on the inter-particle contact number than on contact ratio (contact area).

514



515

516 Fig. 15 Average contact radius ratio and average coordination number increase with the
 517 increment of the mean value of sphericity and roundness. (a) Average contact radius ratio and
 518 (b) Average coordination number

519 6 Conclusion

520 This study develops a method to calculate three-dimensional (3D) sphericity and roundness
 521 of individual particles in a specimen while preserving its internal structure. By comparing the
 522 existing definitions of 3D sphericity, a combination of sphericity and roundness is used to
 523 characterise particle shape. Sphericity and roundness are required to be calculated on a smooth
 524 particle surface which is achieved using a proposed framework including CT techniques,
 525 imaging techniques and the Tabuin smoothing method. The reliability of the smoothing
 526 framework is presented by its application to voxel-constructed sphere, ellipsoid, Ottawa sand
 527 particle and Angular sand particle.

528 By calculating the sphericity and roundness of individual particles in five natural sands, we
 529 found that irregular particles tend to show lower sphericity and roundness than more regular
 530 (and round) particles, and granular packings formed by irregular particles show a wider range
 531 of sphericity and roundness. Moreover, after estimating the thermal conductivity of the natural
 532 sands, we observed that granular materials with a higher average value of sphericity and
 533 roundness show a tendency to boost higher thermal conductivity. This is because we also found
 534 that lower average value of sphericity and roundness may lead to lower average coordination
 535 number and contact radius ratio, important parameters governing heat transfer at the particle
 536 scale. Granular assemblies containing more platy particles may exhibit higher average contact
 537 radius ratio compared to other assemblies with the same average coordination number. The five
 538 materials in this paper have different friction coefficients, the effect of the friction coefficient
 539 on the coordination number and contact ratio requires to be quantified in future work.

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